

Gravitinos can not exist in N-dimensional de Sitter space unless N=4



Based on https://arxiv.org/abs/2206.09851 [1]



1. A few words about dS_N

Einstein equations for N-dimensional de Sitter spacetime (dS_N) :

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + \Lambda g_{\mu\nu} = 0,$$

where $\Lambda = [(N-2)(N-1)]/2\mathcal{R}^2$. (\mathcal{R} is the de Sitter radius. Below $\mathcal{R} = 1$.)

- Topology of dS_N : $\mathbb{R} \times S^{N-1}$. The de Sitter (dS) group is SO(N,1).
- Elementary particles in $dS_N \to \text{unitary}$ irreducible representations (reps.) of the dS algebra spin(N,1).

2. Background and main aim

• The gravitino field (i.e. massless spin-3/2 field) is described by a vector-spinor Ψ_{μ}

$$\gamma^a \nabla_a \Psi_\mu = -M \Psi_\mu$$

$$\gamma^a \Psi_a = \nabla^a \Psi_a = 0,$$
(1)

with (imaginary!) mass parameter M = i(N - 2)/2 [2].

• Set of 'physical' mode solutions of (1): $\{\psi_{\mu}^{(\sigma)}\}$. They form a rep. of spin(N,1). (σ represents quantum numbers.)

These reps. have not been studied in detail.

- Our aim: 1) Construct the modes and 2) study their group-theoretic properties (for $N \ge 3$).
- Extra motivation: All known supergravity (SUGRA) theories in de Sitter are non-unitary. Our analysis for the free gravitino is relevant to linearised SUGRA \rightarrow might be the first step to shed light on the problem.

3. Gauge-invariance of Ψ_{μ}

- Ψ_{μ} is massless (i.e. gauge-invariant).
- Although imaginary M is counter-intuitive, it is predicted by rep. theory [2].
- How does gauge-invariance manifest itself? Eq. (1) admits 'pure gauge' solutions

$$\psi_{\mu}^{(PG)(\sigma)} = (\nabla_{\mu} + \frac{i}{2}\gamma_{\mu})\epsilon^{(\sigma)}$$

 $(\epsilon^{(\sigma)})$ is a spinor). If an invariant norm exists, then it vanishes for $\psi_{\mu}^{(PG)(\sigma)}$.

4. Group-theoretic tools

- Killing vectors (KV's) of dS_N are generators of spin(N, 1).
- Transformation of $\psi_{\mu}^{(\sigma)}$ generated by a KV $\xi^{\mu} \to \text{Lie-Lorentz derivative}$: $\mathscr{L}_{\xi}\psi_{\mu}^{(\sigma)}$ [3].
- $\mathscr{L}_{\xi}\psi_{\mu}^{(\sigma)}$ can be expressed as a linear combination $\mathscr{L}_{\xi}\psi_{\mu}^{(\sigma)} = \sum_{\sigma'} a_{\sigma'} \psi_{\mu}^{(\sigma')}$.
- Unitarity ensures positivity of probabilities. The rep. is unitary if: i) dS invariant scalar product $\langle \psi^{(\sigma)}, \psi^{(\sigma'')} \rangle$

$$\langle \mathscr{L}_{\xi} \psi^{(\sigma)}, \psi^{(\sigma'')} \rangle + \langle \psi^{(\sigma)}, \mathscr{L}_{\xi} \psi^{(\sigma'')} \rangle = 0$$
 (2)

for all KV's ξ and for all σ, σ'' and

ii) the scalar product is positive-definite.

5. Method

- We investigate the unitarity of the reps. The steps of our method are:
- 1) Construct the modes explicitly.
- 2) Choose the KV ξ^{μ} to be a specific dS boost.
- 3) Calculate the coefficients $a_{\sigma'}$ in $\mathcal{L}_{\xi}\psi_{\mu}^{(\sigma)} = \sum_{\sigma'} a_{\sigma'} \psi_{\mu}^{(\sigma')}$.
- 4) Plug $\mathscr{L}_{\xi}\psi_{\mu}^{(\sigma)} = \sum_{\sigma'} a_{\sigma'} \psi_{\mu}^{(\sigma')}$ into the dS invariance condition (2) \rightarrow find conditions for the norms of modes.

6. Main results

We find for the spin(N, 1) reps.:

- For N odd: There do not exist dS invariant scalar products that do not vanish identically \rightarrow the rep. is not unitary.
- For even N > 4: All dS invariant scalar products are indefinite \rightarrow the rep. is not unitary.
- For N=4: Positive-definite and dS invariant scalar products exist \rightarrow the rep. is unitary.
- CONCLUSION: A unitary theory for the gravitino does not exist on dS_N unless N=4.

References

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- [3] T. Ortín, "A note on Lie-Lorentz derivatives," Classical and Quantum Gravity, vol. 19, no. 15, pp. L143–L149, 2002.